

**Reflection and transmission of line-source excited pulsed EM fields at a thin, high-contrast layer with dielectric and conductive properties**

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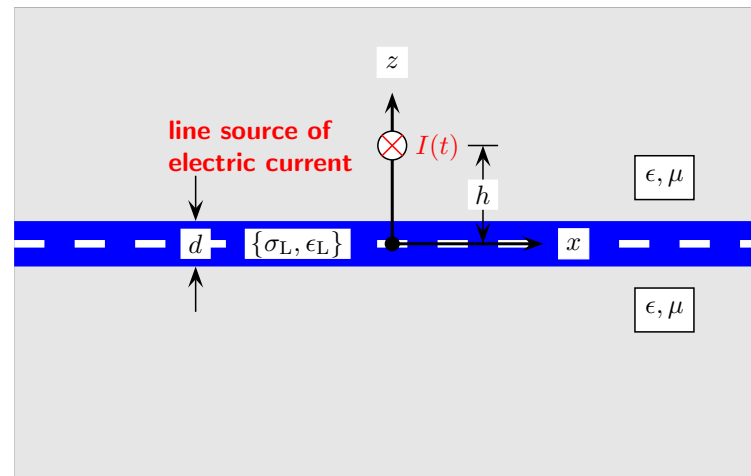
**SYNOPSIS:**

- Problem formulation
- Method of (analytical) solution ( Generalized Cagniard-DeHoop method )
- Illustrative numerical results
- Application to signal integrity analysis in pulsed-field EM radiative inter- and intra-device signal transfer

**AIM OF THE RESEARCH:**

- Studying methods for characterizing pulsed-field EM radiative signal transfer in inter- and intra-microelectronic circuits/devices through the time-domain analysis of appropriate model configurations
- Analyzing the signal integrity associated with the signal transfer in such configurations
- Developing the basis for appropriate time-domain EM Interference regulations pertaining to pulsed-field inter- and intra-microelectronic circuits/devices signal transfer
- Constructing analytical results for the benchmarking of the relevant EM computational procedures

**MODEL CONFIGURATION:** • Line-source of electric current excited thin, high-contrast layer with dielectric and conductive properties



- $G_L = \lim_{d \downarrow 0} \int_{z=-d/2}^{d/2} \sigma_L(z) dz$  (thin-sheet conductance)
- $C_L = \lim_{d \downarrow 0} \int_{z=-d/2}^{d/2} \epsilon_L(z) dz$  (thin-sheet capacitance)

**2D MAXWELL FIELD EQUATIONS:** •  $\{H_x, E_y, H_z, J_y\}(x, z, t); \partial_y = 0$ 

$$\bullet \begin{bmatrix} 0 & -\partial_z & 0 \\ \partial_z & 0 & -\partial_x \\ 0 & \partial_x & 0 \end{bmatrix} \begin{bmatrix} H_x \\ 0 \\ H_z \end{bmatrix} - \epsilon \partial_t \begin{bmatrix} 0 \\ E_y \\ 0 \end{bmatrix} = \begin{bmatrix} 0 \\ J_y \\ 0 \end{bmatrix}$$

$$\bullet \begin{bmatrix} 0 & -\partial_z & 0 \\ \partial_z & 0 & -\partial_x \\ 0 & \partial_x & 0 \end{bmatrix} \begin{bmatrix} 0 \\ E_y \\ 0 \end{bmatrix} + \mu \partial_t \begin{bmatrix} H_x \\ 0 \\ H_z \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

**THIN-SHEET BOUNDARY CONDITIONS:**

$$\bullet E_y(x, z, t) \Big|_{z=-d/2}^{z=d/2} = 0 \quad \bullet H_x(x, z, t) \Big|_{z=-d/2}^{z=d/2} = (G_L + C_L \partial_t) E_y(x, 0, t)$$

## FIELD EXCITATION (ELECTRIC-CURRENT LINE SOURCE):

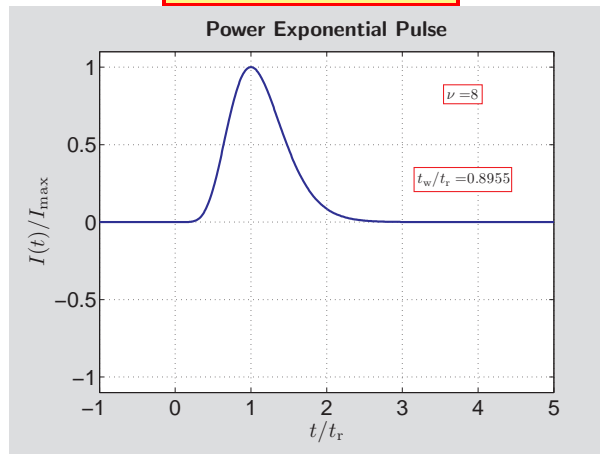
- $J_y(x, z, t) = I(t)\delta(0, h)$
- $I(t) =$  source signature

## POWER EXPONENTIAL PULSE (3-PARAMETER UNIPOLAR PULSE):

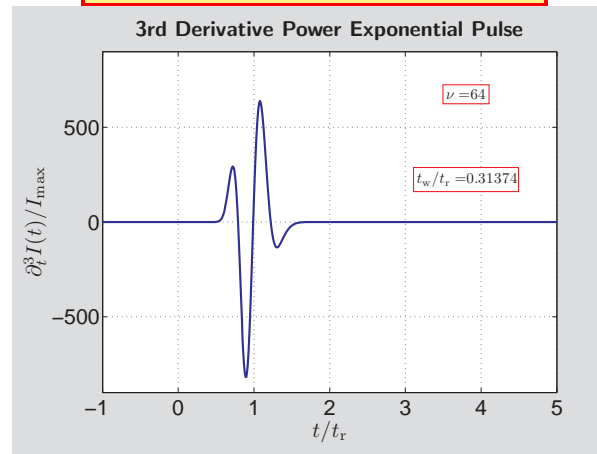
$$I(t) = I_{\max} \left( t/t_r \right)^{\nu} \exp \left[ -\nu \left( t/t_r - 1 \right) \right] H(t)$$

- $I_{\max}$  = pulse amplitude
- $t_r$  = pulse rise time
- $\nu$  = pulse rising power

### Model pulse



### Typical digital pulse



- **TOTAL FIELD**  $\{H_x, E_y, H_z\} =$
- **INCIDENT FIELD**  $\{H_x^i, E_y^i, H_z^i\} +$
- **SCATTERED FIELD**  $\{H_x^s, E_y^s, H_z^s\}$

### CAGNIARD-DEHOOP METHOD

• Time Laplace transformation ( $\hat{\cdot}$ )	$\hat{\partial}_t = s$	$s \in \mathbb{R}, s > 0$
• Waveslowness representation ( $\tilde{\cdot}$ )	$\tilde{\partial}_x = -i\alpha s$	$\alpha \in \mathbb{R}$
• Cagniard-DeHoop representation	$i\alpha = p$	$p \in \mathbb{C}$
<b>waveslowness domain</b> $\{p, z, s\}$ $\mapsto$		<b>spacetime domain</b> $\{x, z, \tau\}$
• Cagniard-DeHoop contour deformation	$p \rightarrow \bar{p}(x, z, \tau)$	$\tau = \text{time} \in \mathbb{R}$

### • INCIDENT FIELD:

- $\tilde{E}_y^i(p, z, s) = \mu s \hat{I}(s) \tilde{G}^i(p, z, s)$
- $\tilde{G}^i(p, z, s) = \frac{\exp[-s\gamma(p)|z - h|]}{2 s \gamma(p)}$ 
  - $\gamma(p) = (1/c^2 - p^2)^{1/2} \quad c = (\epsilon\mu)^{-1/2} \quad \text{Re}[\gamma(p)] \geq 0$   
(slowness-domain propagation coefficient)

### • SCATTERED FIELD:

- $\tilde{E}_y^s(p, z, s) = \mu s \hat{I}(s) \tilde{S}_L(p, s) \tilde{G}^s(p, z, s)$
- $\tilde{G}^s(p, z, s) = \frac{\exp[-s\gamma(p)(h + |z|)]}{2 s \gamma(p)}$ 
  - $\tilde{S}_L(p, s)$  = slowness-domain scattering coefficient layer



**INCIDENT FIELD:**

- $\hat{E}_y^i(x, z, s) = \mu s \hat{I}(s) \frac{1}{2\pi i} \int_{p=-i\infty}^{i\infty} \frac{\exp\{-s[p x + \gamma(p)|z - h|]\}}{2 s \gamma(p)} dp$
- $p \rightarrow p^i(x, z, \tau) : p^i x + \gamma(p^i)|z - h| = \tau \mid \bullet T^i \leq \tau < \infty$  (CdH-path)
- $\hat{E}_y^i(x, z, s) = \mu s \hat{I}(s) \int_{\tau=T^i}^{\infty} \exp(-s\tau) \frac{1}{2\pi(\tau^2 - T^{i2})^{1/2}} d\tau$
- $T^i = (x^2 + |z - h|^2)^{1/2}/c$  (arrival time)

**SCATTERED FIELD:**

- $\hat{E}_y^s(p, z, s) = \mu s \hat{I}(s) \frac{1}{2\pi i} \int_{p=-i\infty}^{i\infty} S_L(p, s) \frac{\exp\{-s[p x + \gamma(p)(h + |z|)]\}}{2 s \gamma(p)} dp$
- $p \rightarrow p^s(x, z, \tau) : p^s x + \gamma(p^s)(h + |z|) = \tau \mid \bullet T^s \leq \tau < \infty$  (CdH-path)
- $\hat{E}_y^s(p, z, s) = \mu s \hat{I}(s) \int_{\tau=T^s}^{\infty} \exp(-s\tau) \text{Re}[\tilde{S}_L(p^s, s)] \frac{1}{2\pi(\tau^2 - T^{s2})^{1/2}} d\tau$
- $T^s = (x^2 + (h + |z|)^2)^{1/2}/c$  (arrival time)

**INCIDENT FIELD:**

$$\bullet E_y^i(x, z, t) = \mu \partial_t I(t) \overset{(t)}{*} \left[ \frac{1}{2\pi(\tau^2 - T^i)^{1/2}} \right] H(t - T^i)$$

**SCATTERED FIELD:**

$$\bullet E_y^s(x, z, t) = \mu \partial_t I(t) \overset{(t)}{*} \left[ - \frac{1}{2\pi(t^2 - T^s)^{1/2}} + \int_{\tau=T^s}^t K_L(\tau, t - \tau) \frac{1}{2\pi(\tau^2 - T^s)^{1/2}} d\tau \right] H(t - T^s)$$

$$\bullet K_L(\tau, t) = \text{Re} \left[ \beta_L(p^s) \exp\{-[\beta_L(p^s) + G_L/C_L]t\} \right] H(t)$$

$$\bullet \beta_L(p^s) = \frac{2\gamma(p^s)}{\mu C_L}$$

$\overset{(t)}{*}$  = time convolution

Field expressions (time domain)

**ILLUSTRATIVE NUMERICAL  
RESULTS**

## CONCLUSION:

- Closed-form analytical field expressions facilitate quantifying the sensitivity of the pulse shape of the scattered field to the different configurational parameters
- The constructed time-domain algorithm can serve as a tool towards the study of the signal integrity related to the source/receiver signal transfer in configurations with thin, highly contrasting layers with conductive and capacitive properties
- Further research  $\implies$  extension to the corresponding 3D (point-source excited) model, with particular application to inter- and intra-device radiative transfer of digital signals