

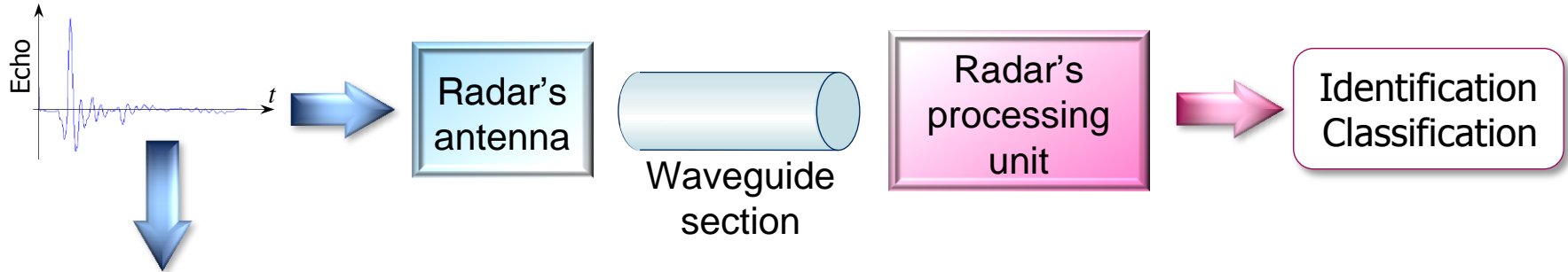
Pulse shape distortion in closed-waveguide axial modal signal transfer an analytic time-domain study

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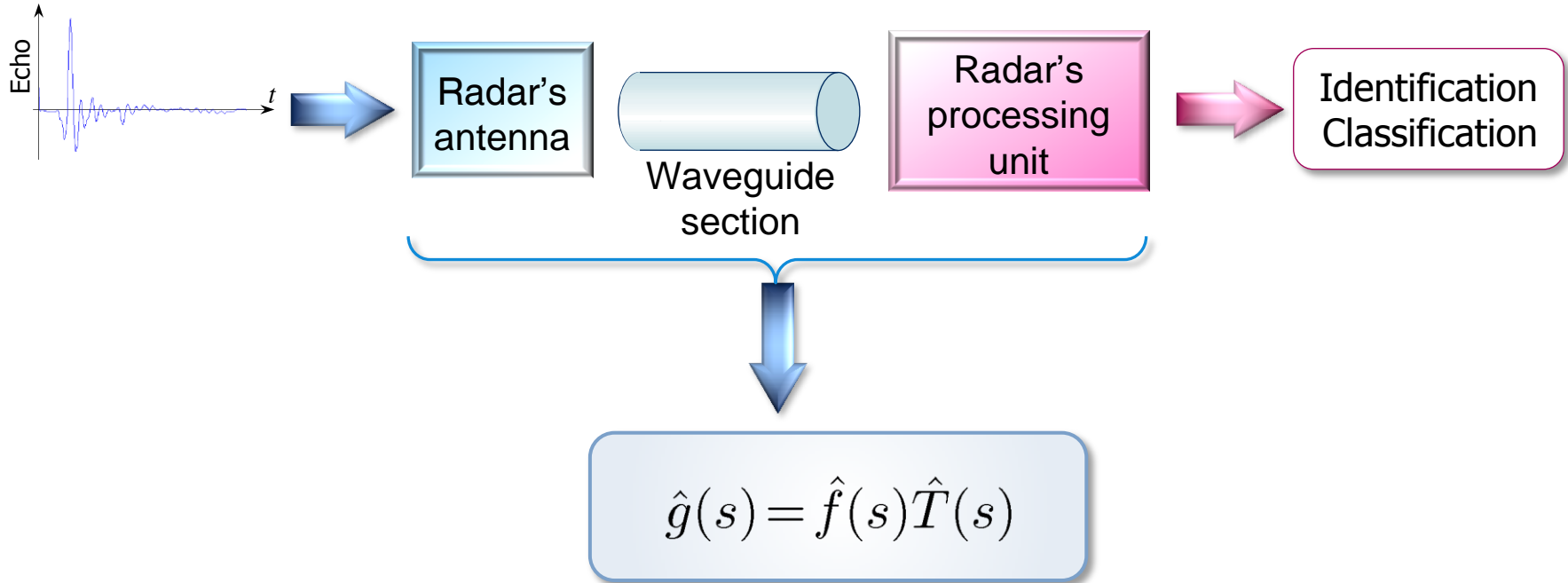
Pulsed radar scenario



$$f(t)H(t) \xrightarrow{\text{LT}} \hat{f}(s) = \int_{t=0}^{\infty} f(t) \exp(-st) dt$$
$$s \in \mathbb{C}, \operatorname{Re}(s) > 0$$

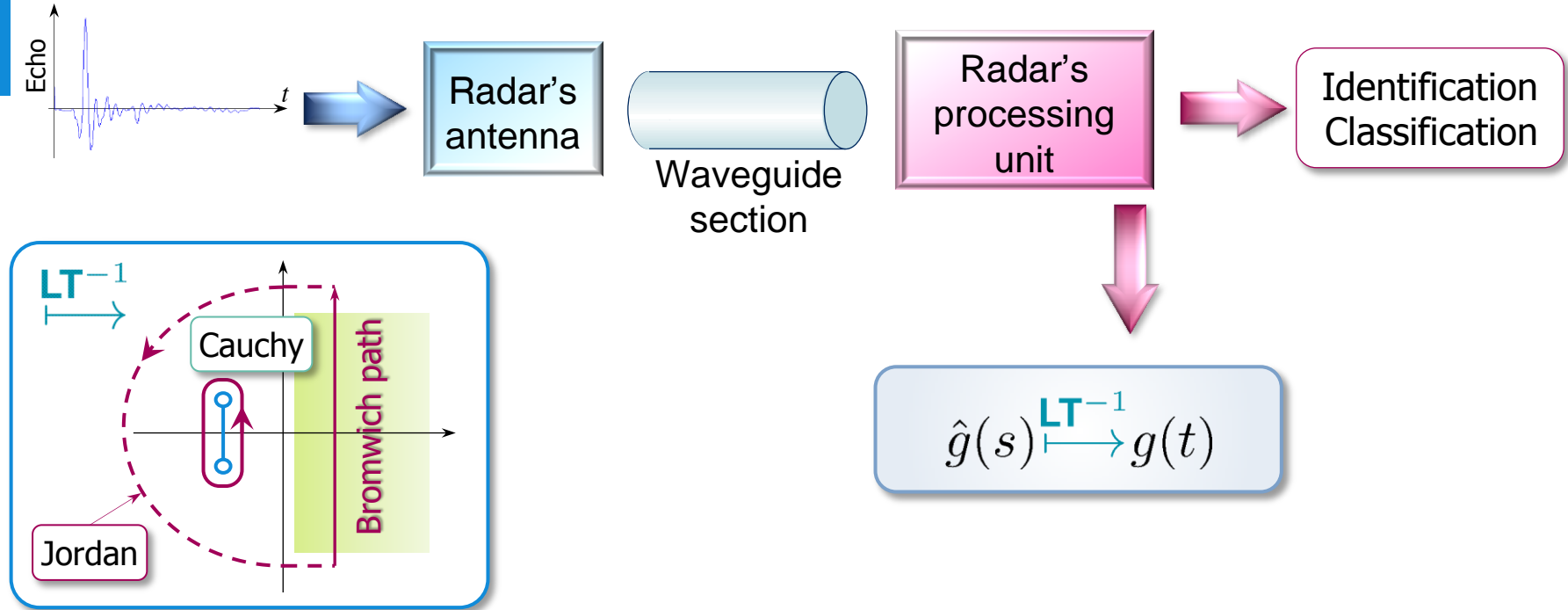
- Causal echo \Rightarrow via the **time Laplace transform**

Pulsed radar scenario



- Linear, time-invariant, passive systems 
causal response

Pulsed radar scenario



- The dispersion properties of the transfer path \rightarrow **determinant for the detection ability**

Synopsis

- Propagation / propagators
- **Time-domain** modal expansion in closed waveguides
- Illustrative numerical experiments
- Conclusions

Prerequisites

- **Material parameters:**

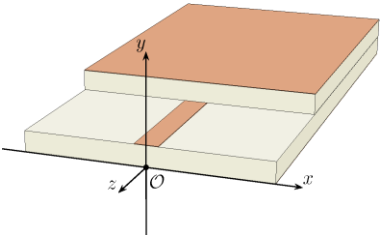
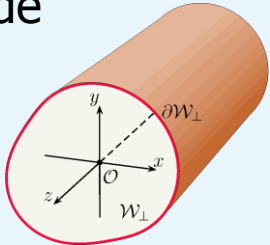
- $\varepsilon > 0$ = electric permittivity
- $\sigma \geq 0$ = electric conductivity
- $\mu > 0$ = magnetic permeability
- $\kappa \geq 0$ = linear magnetic loss coefficient
- $c = (\varepsilon\mu)^{-1/2}$ = wave speed

- **Instrument:** the time Laplace transform $f(t)H(t) \xrightarrow{\text{LT}} \hat{f}(s)$

- **Impedances:**

- $\hat{\zeta} = \mu(s + \kappa/\mu)$ = longitudinal impedance of the medium
- $\hat{\eta} = \varepsilon(s + \sigma/\varepsilon)$ = transverse admittance of the medium

Propagation / propagators

Propagation path	Propagator
Free space	$\frac{\exp(-\gamma R)}{4\pi R}$ <p>$R =$ radial distance $\gamma = s/c_0 =$ propagation coeff.</p>
Stripline 	$\exp(\mp\gamma z)$ <p>$z =$ axial distance $\gamma = \frac{1}{c} [(s + \alpha)(s + \beta)]^{1/2}$ Material dispersion Boltzmann relaxation coefficients $\alpha \geq 0, \beta \geq 0$</p>
Closed waveguide 	$\exp(\mp\gamma_n z)$ <p>$\gamma_n = (\gamma^2 + \kappa_n^2)^{1/2}$ Modal geometrical dispersion modal eigenvalue</p>

Closed waveguides

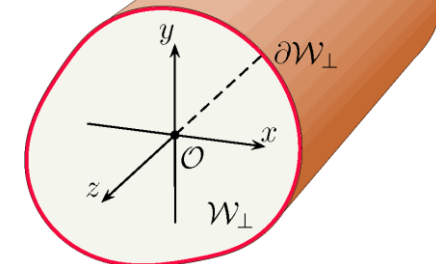
- Time Laplace transformed Maxwell eqs.

$$\begin{aligned}\partial \times \hat{H} - \hat{\eta} \hat{E} &= \hat{J}, \\ \partial \times \hat{E} + \hat{\zeta} \hat{H} &= -\hat{K},\end{aligned}$$

- ∂ = spatial differential operator
- \hat{E} = electric field strength
- \hat{H} = magnetic field strength
- \hat{J} = volume source density of electric current
- \hat{K} = volume source density of magnetic current

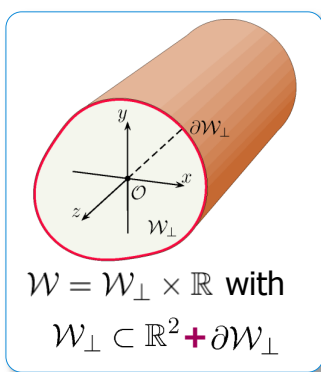
- Geometric vectorial decomposition

- \parallel = parallel to waveguide axis
- \perp = perpendicular to waveguide axis



$$\begin{aligned}\mathcal{W} &= \mathcal{W}_{\perp} \times \mathbb{R} \text{ with} \\ \mathcal{W}_{\perp} &\subset \mathbb{R}^2 + \partial \mathcal{W}_{\perp}\end{aligned}$$

Closed waveguides



- Modal decomposition \Rightarrow modal constituents:
 - transverse magnetic (TM) modes \Rightarrow in terms of E_{\parallel}
 - transverse electric (TE) modes \Rightarrow in terms of H_{\parallel}

- Time Laplace-transformed scalar wave equations

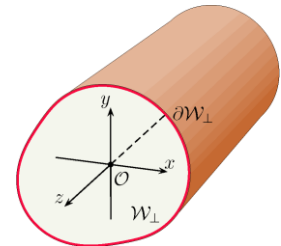
$$\left(\partial_{\perp} \cdot \partial_{\perp} + \partial_{\parallel} \partial_{\parallel} - \hat{\eta} \hat{\zeta} \right) \left\{ \hat{E}_{\parallel}, \hat{H}_{\parallel} \right\} = - \left\{ \hat{Q}_{\parallel}^E, \hat{Q}_{\parallel}^H \right\}$$

- Time Laplace-transformed Green's function of the dissipative scalar wave equation

$$\left[\partial_{\perp} \cdot \partial_{\perp} + \partial_{\parallel} \partial_{\parallel} - c^{-2} (s + \alpha) (s + \beta) \right] \hat{G}$$

$$= -\delta \left(\mathbf{x}_{\perp}, \mathbf{x}'_{\perp}, x_{\parallel} - x'_{\parallel} \right)$$

Modal decomposition



$\mathcal{W} = \mathcal{W}_\perp \times \mathbb{R}$ with
 $\mathcal{W}_\perp \subset \mathbb{R}^2 + \partial\mathcal{W}_\perp$

- Modal expansion functions \Rightarrow solutions of the transverse Helmholtz equation

$$(\partial_{\parallel}\partial_{\parallel} + \kappa_n^2) \varphi_n = 0, \text{ for } \mathbf{x}_\perp \in \mathcal{W}_\perp$$

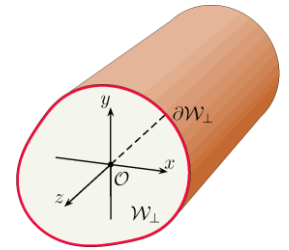
- Dirichlet B.C.-s on \mathcal{W}_\perp \Rightarrow TM-modes
- Neumann B.C.-s on \mathcal{W}_\perp \Rightarrow TE-modes
- normalisation:

$$\delta(\mathbf{x}_\perp, \mathbf{x}'_\perp) = \sum_n N_n^{-1} \varphi_n(\mathbf{x}_\perp) \varphi_n(\mathbf{x}'_\perp)$$

with

$$N_n = \int_{\mathcal{W}_\perp} \varphi_n(\mathbf{x}_\perp) \varphi_n(\mathbf{x}_\perp) dA$$

Modal decomposition



$\mathcal{W} = \mathcal{W}_\perp \times \mathbb{R}$ with
 $\mathcal{W}_\perp \subset \mathbb{R}^2 + \partial\mathcal{W}_\perp$

- Spectral Green's function

$$\hat{G} = \sum_n N_n^{-1} \varphi_n(\mathbf{x}_\perp) \varphi_n(\mathbf{x}'_\perp) \hat{g}_n(\mathbf{x}_\parallel - \mathbf{x}'_\parallel, s)$$

in which

$$(\partial_\parallel \partial_\parallel - \gamma_n^2) \hat{g}_n(\mathbf{x}_\parallel - \mathbf{x}'_\parallel, s) = -\delta(\mathbf{x}_\parallel - \mathbf{x}'_\parallel)$$

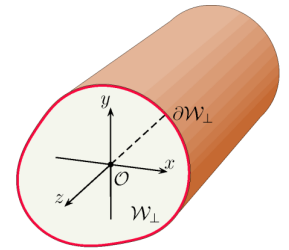
- Modal characteristic

$$\gamma_n^2 = \frac{(s + \alpha)(s + \beta)}{c^2} + \kappa_n^2 = \frac{[s + (\alpha + \beta)/2]^2}{c^2} + \Omega_n^2$$

with

$$\Omega_n^2 = \kappa_n^2 - \frac{c^{-2} (\alpha - \beta)^2}{4}$$

Modal decomposition



$\mathcal{W} = \mathcal{W}_\perp \times \mathbb{R}$ with
 $\mathcal{W}_\perp \subset \mathbb{R}^2 + \partial\mathcal{W}_\perp$

- Time-domain Green's function

$$G_n = \sum_n N_n^{-1} \varphi_n(\mathbf{x}_\perp) \varphi_n(\mathbf{x}'_\perp) g_n(\mathbf{x}_\parallel - \mathbf{x}'_\parallel, t)$$

with

$$g_n(\mathbf{x}_\parallel - \mathbf{x}'_\parallel, t) = \frac{c}{2} \exp\left[-\frac{(\alpha + \beta)t}{2}\right] J_0\left[c\Omega_n \left(t^2 - |\mathbf{x}_\parallel - \mathbf{x}'_\parallel|^2 / c^2\right)^{1/2}\right] \\ H\left[t - |\mathbf{x}_\parallel - \mathbf{x}'_\parallel| / c\right] \quad \text{for } \Omega_n^2 > 0$$

Oscillatory behaviour

Modal decomposition

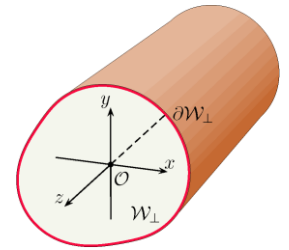
- Time-domain Green's function

$$G_n = \sum_n N_n^{-1} \varphi_n(\mathbf{x}_\perp) \varphi_n(\mathbf{x}'_\perp) g_n(\mathbf{x}_\parallel - \mathbf{x}'_\parallel, t)$$

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Monotonic behaviour



$\mathcal{W} = \mathcal{W}_\perp \times \mathbb{R}$ with
 $\mathcal{W}_\perp \subset \mathbb{R}^2 + \partial\mathcal{W}_\perp$

Illustrative numerical experiments

- **Analysed quantity:**

$$P_n \left(\mathbf{x}_{\parallel} - \mathbf{x}'_{\parallel}, t \right) = g_n \left(\mathbf{x}_{\parallel} - \mathbf{x}'_{\parallel}, t \right) \stackrel{(t)}{*} S(t)$$

with $S(t)$ = feeding signal

- **Normalisation:**

- time: $t_{\text{norm}} = t c \Omega_{n,\text{ref}}$, with $\Omega_{n,\text{ref}}^2 > 0$ sufficiently large

- space: $|\mathbf{x}_{\parallel} - \mathbf{x}'_{\parallel}|_{\text{norm}} = |\mathbf{x}_{\parallel} - \mathbf{x}'_{\parallel}| \Omega_{n,\text{ref}}$

- analysed value: $P_n \left(\mathbf{x}_{\parallel} - \mathbf{x}'_{\parallel}, t \right)$ to its peak absolute value

- **Medium inside the waveguide:**

- vacuum, $\{\alpha, \beta\} = \{0, 0\}$

- Rogers RO4003™, $\{\alpha, \beta\} = \{0.0027, 0\}$

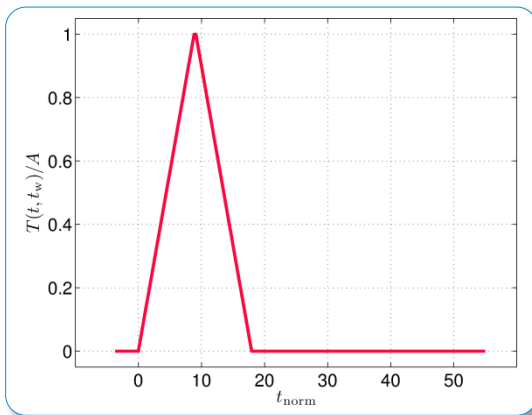
Illustrative numerical experiments

- **Feeding signal:** triangular pulse

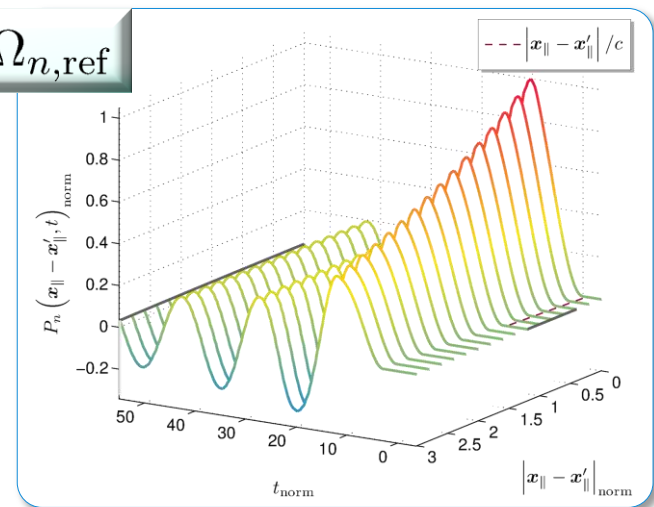
$$T(t, t_w) = A \begin{cases} t/t_w H(t) & \text{for } t \leq t_w \\ 2 - t/t_w & \text{for } t_w < t \leq 2t_w \\ 0 & \text{for } t > 2t_w \end{cases}$$

with $2t_w$ = the pulse width

- **Medium:** vacuum



$$\Omega_n^2 > 0, \quad \Omega_n < \Omega_{n,\text{ref}}$$



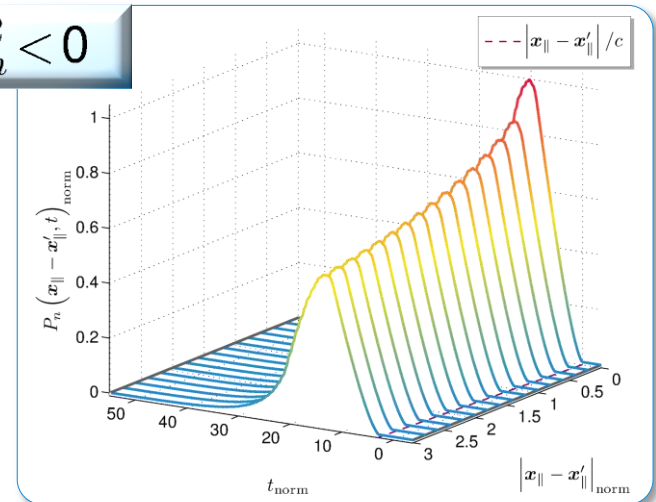
Illustrative numerical experiments

- **Feeding signal:** triangular pulse $\Omega_n^2 < 0$

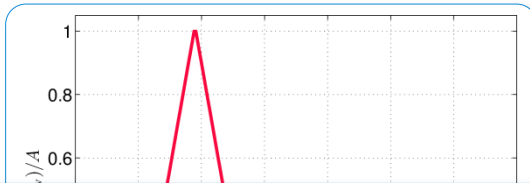
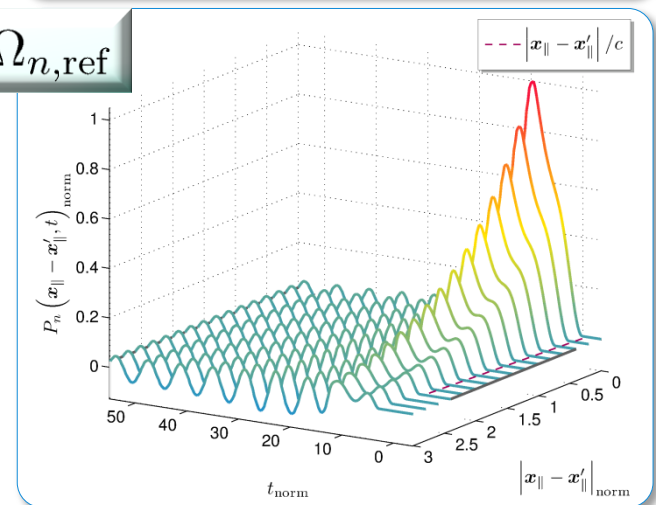
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with $2t_w =$ pulse time width

- **Medium:** Rogers RO4003™



$$\Omega_n^2 > 0, \Omega_n < \Omega_{n,ref}$$



Relevance:

TD pulsed signal transfer \longleftrightarrow FD analyses correspondences

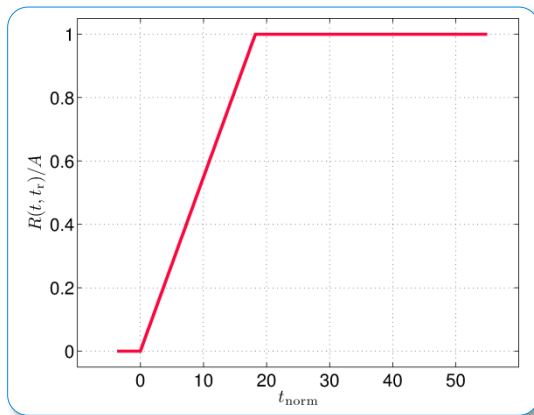
Illustrative numerical experiments

- **Feeding signal:** clipped ramp function

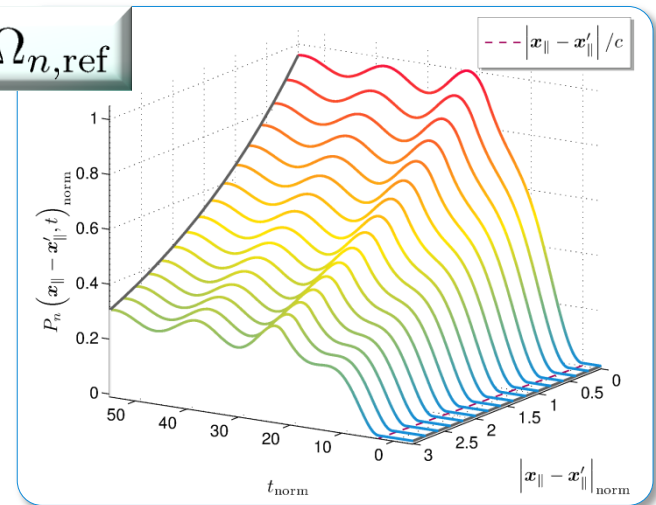
$$R(t, t_r) = A \begin{cases} t/t_r H(t) & \text{for } t \leq t_r \\ 1 & \text{for } t > t_r \end{cases}$$

with $t_r =$ pulse rise time

- **Medium:** vacuum



$$\Omega_n^2 > 0, \quad \Omega_n < \Omega_{n,\text{ref}}$$



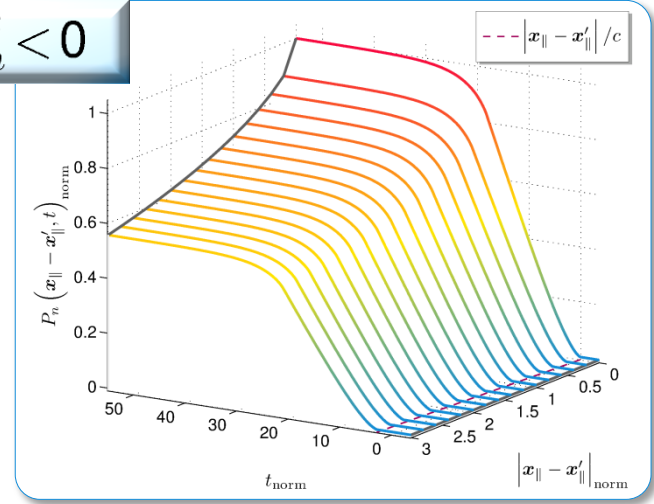
Illustrative numerical experiments

- **Feeding signal:** clipped ramp function $\Omega_n^2 < 0$

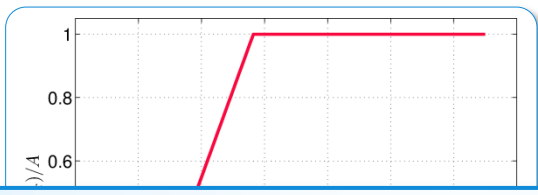
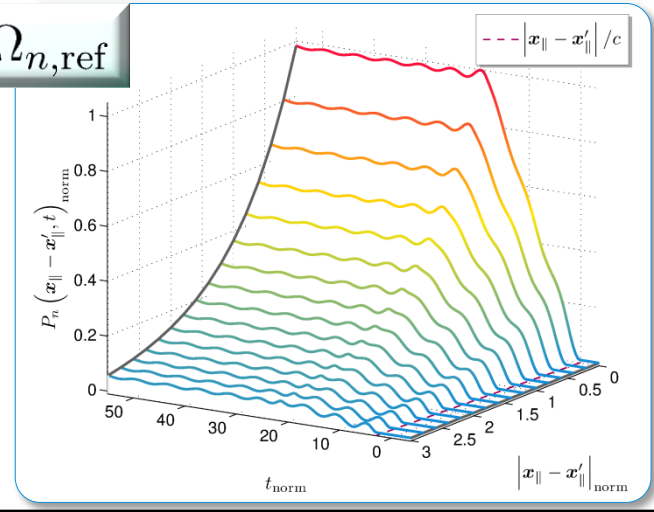
$$R(t, t_r) = A \begin{cases} t/t_r H(t) & \text{for } t \leq t_r \\ 1 & \text{for } t > t_r \end{cases}$$

with $t_r =$ pulse rise time

- **Medium:** Rogers RO4003™




$\Omega_n^2 > 0, \Omega_n < \Omega_{n,ref}$



Relevance:
 TD reconstruction of the pulse shapes of monostatic or multistatic radar returns

Conclusions

- **Examined:** influence of material- and geometry-induced dispersion of a pulsed signal in a closed waveguide
- **General applicability:** effective tools for constructing a deconvolution algorithm that filters out the pertaining dispersion
- **Practical application:** RADAR  reconstruction of a target's dimensions and shape from the irradiation by a field with dedicated pulse time shape